

# Coherent backscattering of bosonic matter-wave in the presence of disorder and interaction

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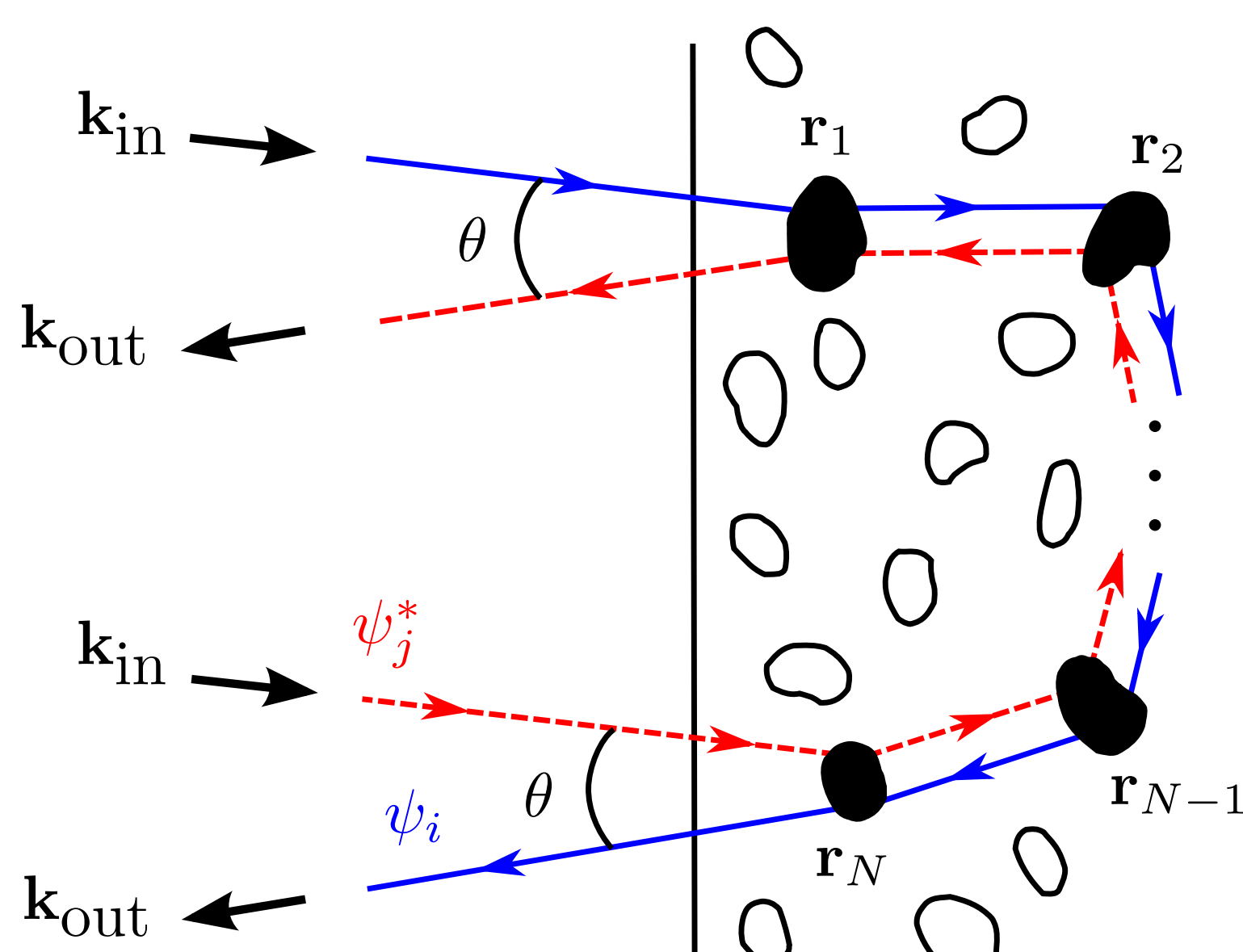


## Abstract

Coherent backscattering, which is an enhancement of the backscattered intensity of a light going through a medium made of point-like scatterers, is known as one of the most robust interference effects. It has been shown, although it is nowadays not fully understood yet, that in the presence of non-linearities this enhancement turns to an inhibition. We propose to study that effect by means of a system in which we study the transport of a Bose-Einstein condensate through Aharonov-Bohm rings in the presence of interaction and disorder. We find that our system is indeed a good candidate to observe the coherent peak's inversion and is also suitable for more feasible theoretical calculations than in the original case.

## Coherent backscattering

Laser light  $I_0$  and wavevector  $k_{in}$  going through a sample composed of point-like scatterers at random positions

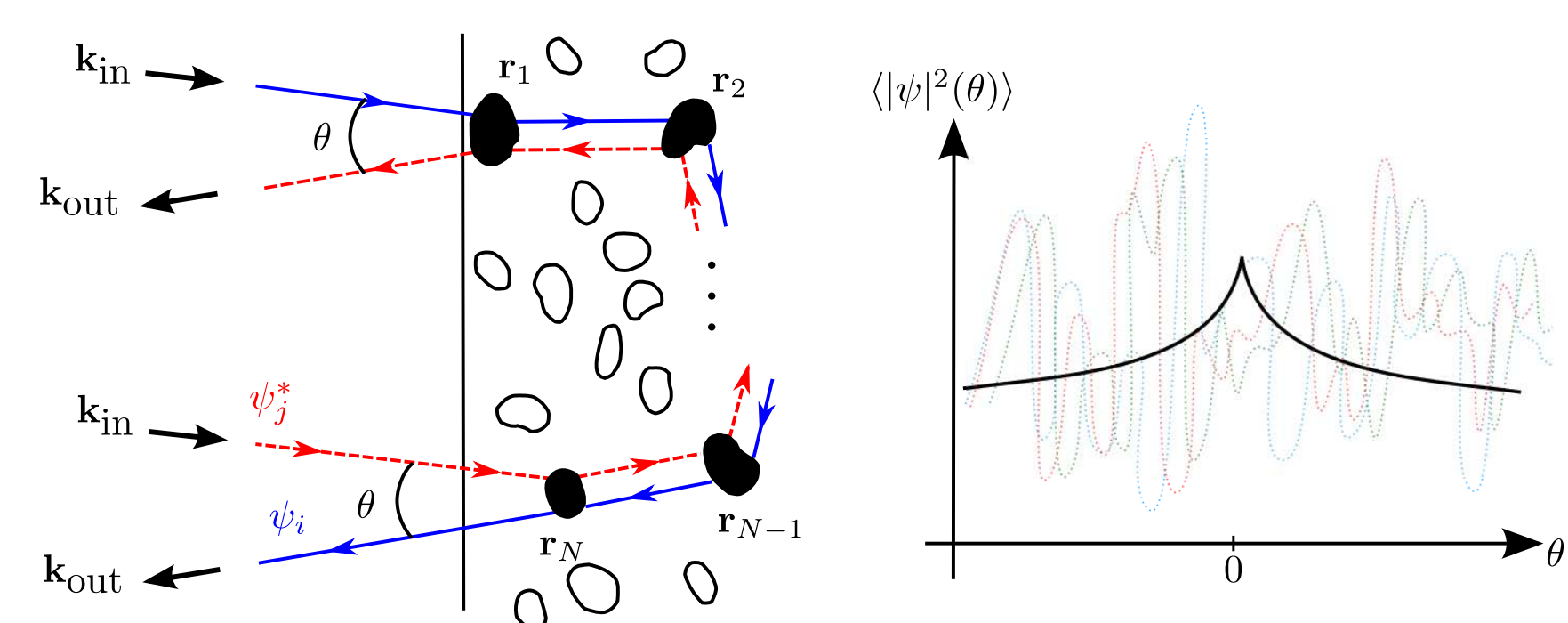


Compute  $\langle I(\theta) \rangle \propto |\psi(\theta)|^2$  with

$$\begin{aligned} \psi &= \sum_{\text{paths } i} \psi_i \Rightarrow |\psi|^2 = \sum_{i,j} \psi_i \psi_j^* \\ &= \sum_{\text{paths } i} |\psi_i|^2 + \sum_{i \neq j} \psi_i \psi_j^* \end{aligned}$$

No interference?  $\Rightarrow$  Ohm's law

Ensemble average over scatterer's positions yields a specific conic pattern

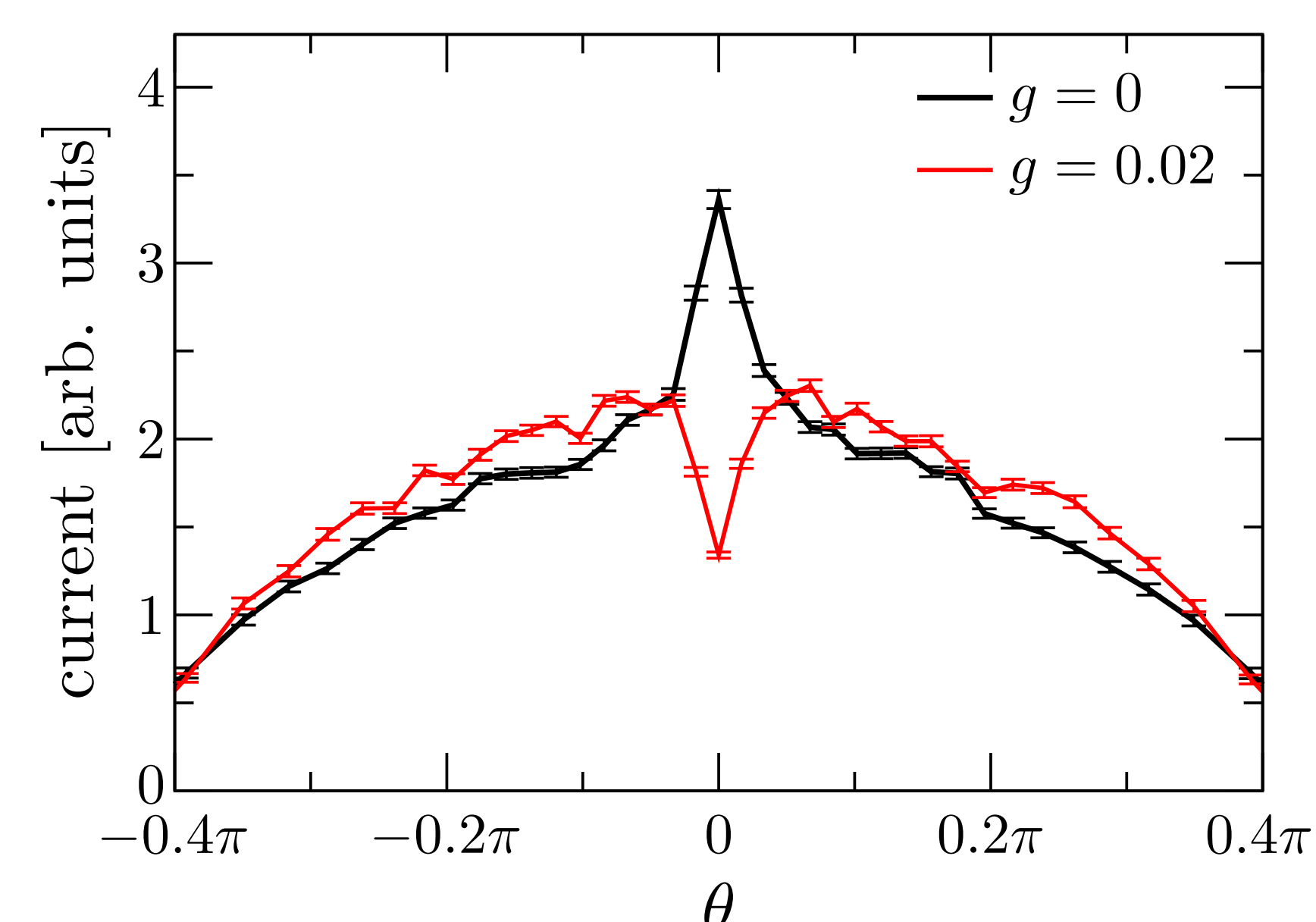


Constructive interferences between time-reversed conjugate paths (same phase acquired due to disorder) around  $\theta = 0$  (same path's length).

## CBS peak inversion

Numerical integration of the Gross-Pitaevskii equation in 2D shows that the interaction strength parameter  $g$  plays an important role

[M. Hartung et al. *Phys. Rev. Lett.* **101**, 020603 (2008)]

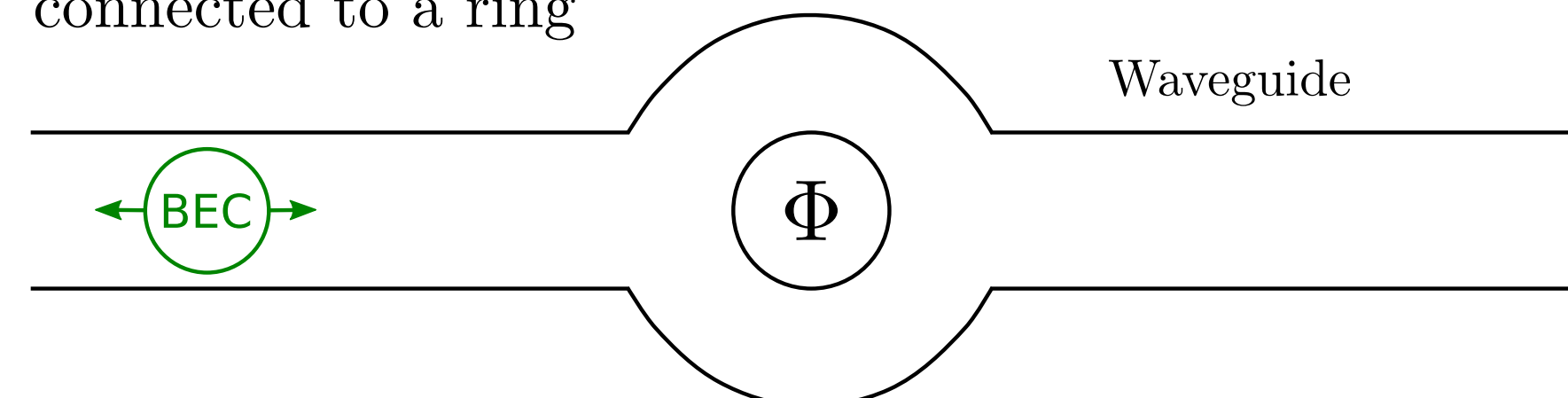


Inverted cone in presence of finite interaction  $\rightarrow$  crossover from constructive to destructive interference

Theoretical calculations beyond the Gross-Pitaevskii approach difficultly feasible in 2D

## Our model

Our system : BEC coupled to 2 semi-infinite waveguides connected to a ring



Usually described by Gross-Pitaevskii equation

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(\mathbf{r}) + V(\mathbf{r})\psi(\mathbf{r}) + g|\psi(\mathbf{r})|^2\psi(\mathbf{r}) = \mu\psi(\mathbf{r})$$

$\rightarrow$  Ok if interaction strength  $g$  "small enough"

Numerical integration of GP equation and truncated Wigner method

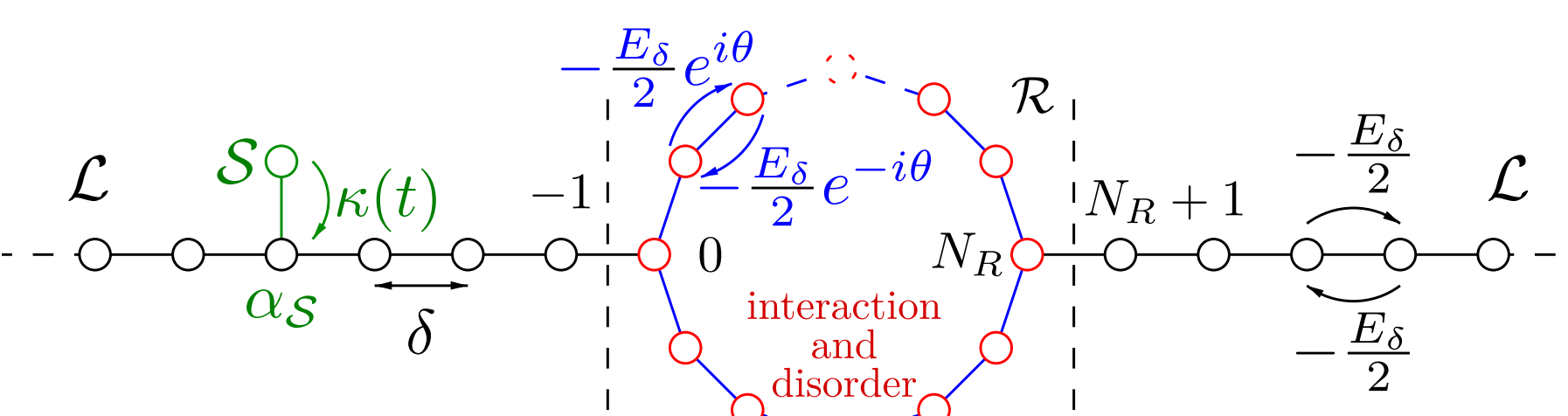
## Theoretical description

Ring geometry connected to two semi-infinite homogeneous leads

Perfect condensation of the reservoir ( $T = 0$  K) with chemical potential  $\mu$

Discretisation of a 1D Bose-Hubbard system

[J. Dujardin et al. *Phys. Rev. A* **91**, 033614 (2015)]



Hamiltonian

$$\hat{H} = \hat{H}_L + \hat{H}_{LR} + \hat{H}_R + \hat{H}_S$$

where

$$\hat{H}_L = \sum_{\alpha \in \mathcal{L}} \left[ E_\delta \hat{a}_\alpha^\dagger \hat{a}_\alpha - \frac{E_\delta}{2} (\hat{a}_{\alpha+1}^\dagger \hat{a}_\alpha + \hat{a}_\alpha^\dagger \hat{a}_{\alpha+1}) \right]$$

$$\hat{H}_{LR} = -\frac{E_\delta}{2} (\hat{a}_{-1}^\dagger \hat{a}_0 + \hat{a}_0^\dagger \hat{a}_{-1} + \hat{a}_{N_R}^\dagger \hat{a}_{N_R+1} + \hat{a}_{N_R+1}^\dagger \hat{a}_{N_R})$$

$$\begin{aligned} \hat{H}_R = & \left[ \sum_{\alpha \in \mathcal{R}} (E_\delta + V_\alpha) \hat{a}_\alpha^\dagger \hat{a}_\alpha - \frac{E_\delta}{2} (\hat{a}_{\alpha-1}^\dagger \hat{a}_\alpha + \hat{a}_{\alpha+1}^\dagger \hat{a}_\alpha) \right. \\ & \left. + g \hat{a}_\alpha^\dagger \hat{a}_\alpha^\dagger \hat{a}_\alpha \hat{a}_\alpha \right] \end{aligned}$$

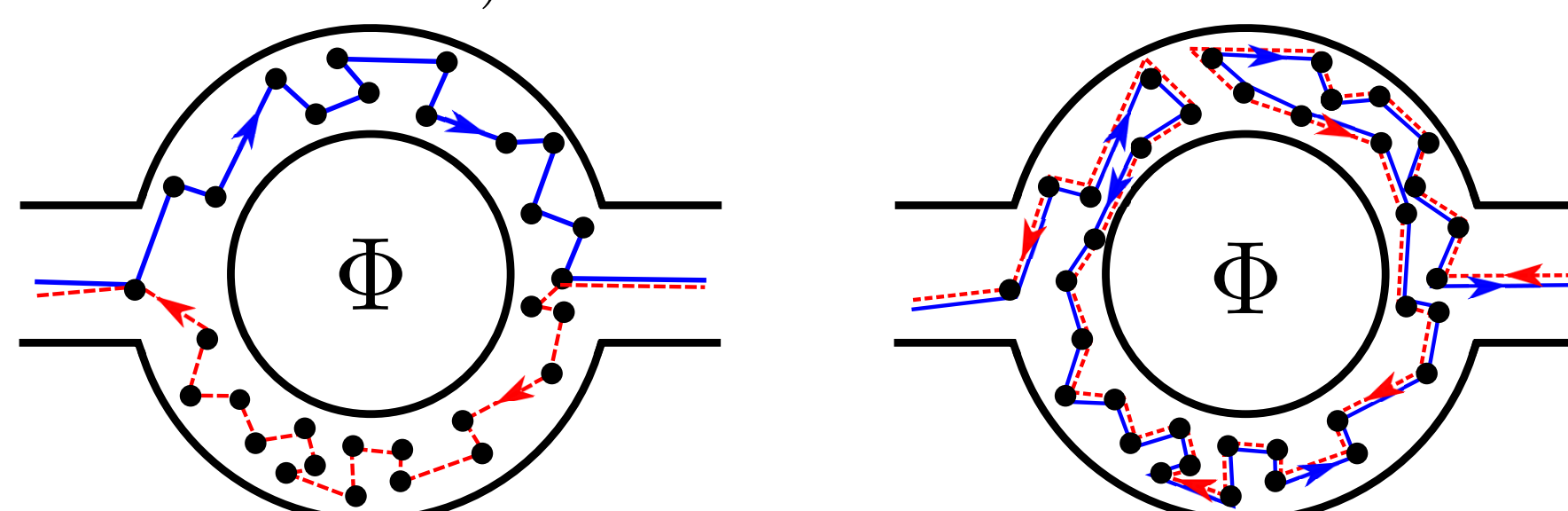
$$\hat{H}_S = \kappa(t) \hat{a}_{\alpha_S}^\dagger \hat{b} + \kappa^*(t) \hat{b}^\dagger \hat{a}_{\alpha_S} + \mu \hat{b}^\dagger \hat{b}$$

with :

- $\hat{a}_\alpha$  ( $\hat{b}$ ) and  $\hat{a}_\alpha^\dagger$  ( $\hat{b}^\dagger$ ) the annihilation and creation operators at site  $\alpha$  (of the source),
- $E_\delta \propto 1/\delta^2$  the on-site energy,
- $V_\alpha$  the disorder potential at site  $\alpha$ ,
- $g$  the interaction strength,
- $N \rightarrow \infty$  the number of Bose-Einstein condensed atoms within the source,
- $\kappa(t) \rightarrow 0$  the coupling strength.

## CBS within a ring

Time-reversed paths are exactly the same (same experienced disorder)



Constructive interference between those paths

Enhanced backscattering probability  $\rightarrow$  Aronov-Altshuler-Spivak oscillations

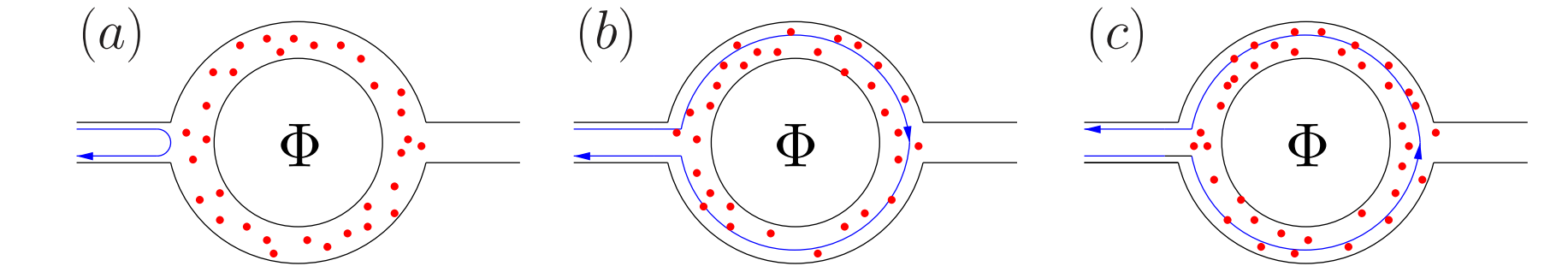
Transmission easier to compute than reflection

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## Higher order interferences

How to highlight the CBS contribution ?

Schematic approach of the problem



[Ihn T., *Semiconductor nanostructures*, Oxford (2010)]

The reflection probability is given by

$$\mathcal{R} = |r_0 + r_1 e^{i\Phi} + r_1 e^{-i\Phi} + \dots|^2 = |r_0|^2 + |r_1|^2 + \dots \quad (1)$$

$$+ 4|r_0| \cdot |r_1| \cos \Lambda \cos \Phi + \dots \quad (2)$$

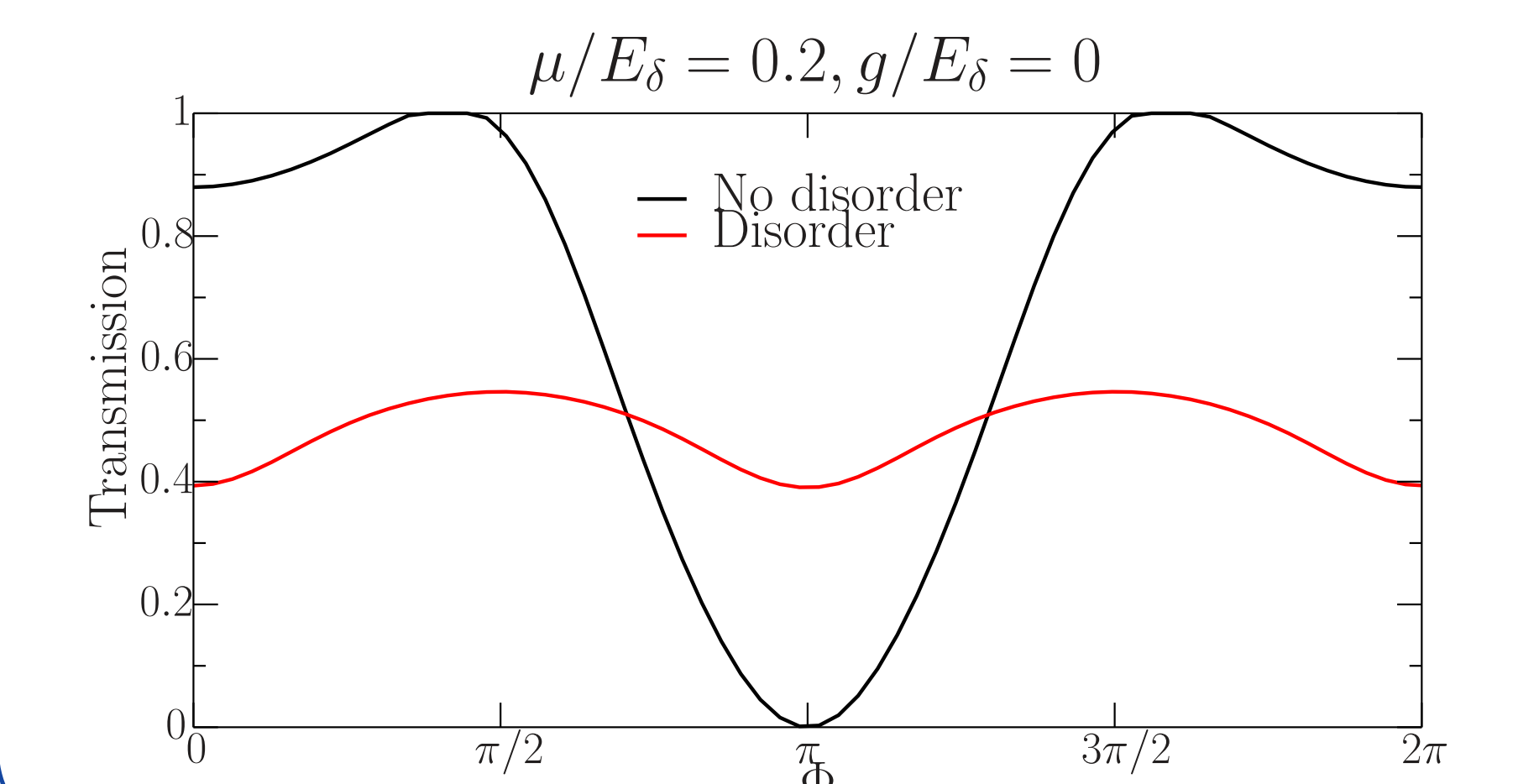
$$+ 2|r_1|^2 \cos(2\Phi) + \dots \quad (3)$$

with  $\Lambda$  the disorder-dependent phase accumulated after one turn with  $\Phi = 0$ .

- (1) no  $\Phi$ -dependence, classical contributions
- (2)  $\Phi$ -periodicity, AB contribution, damped to zero when averaged over the disorder
- (3)  $\Phi/2$ -periodicity, AAS contribution, robust to average over the disorder

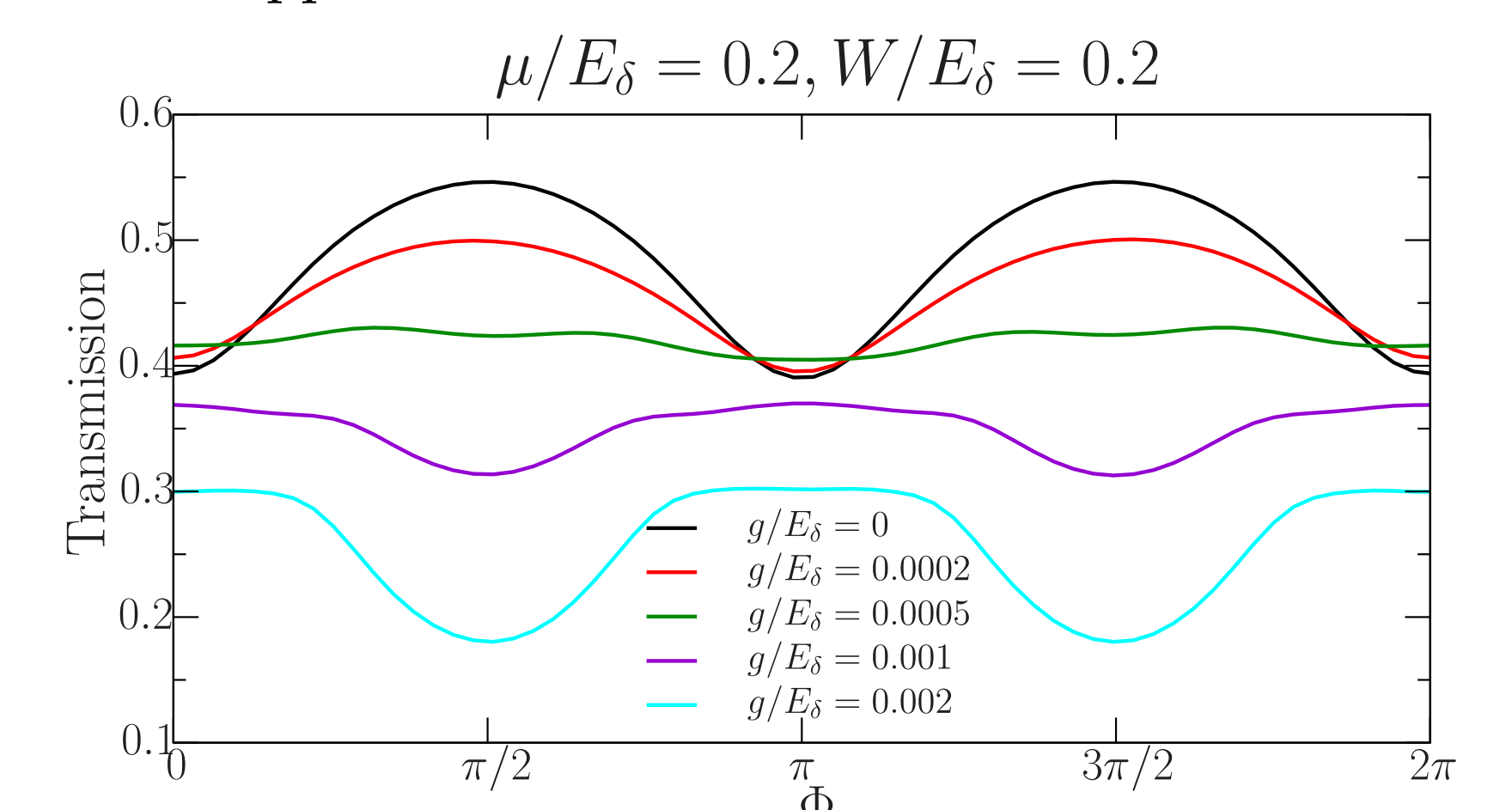
Ensemble average over disorder

$\rightarrow$  Appearance of AAS oscillations

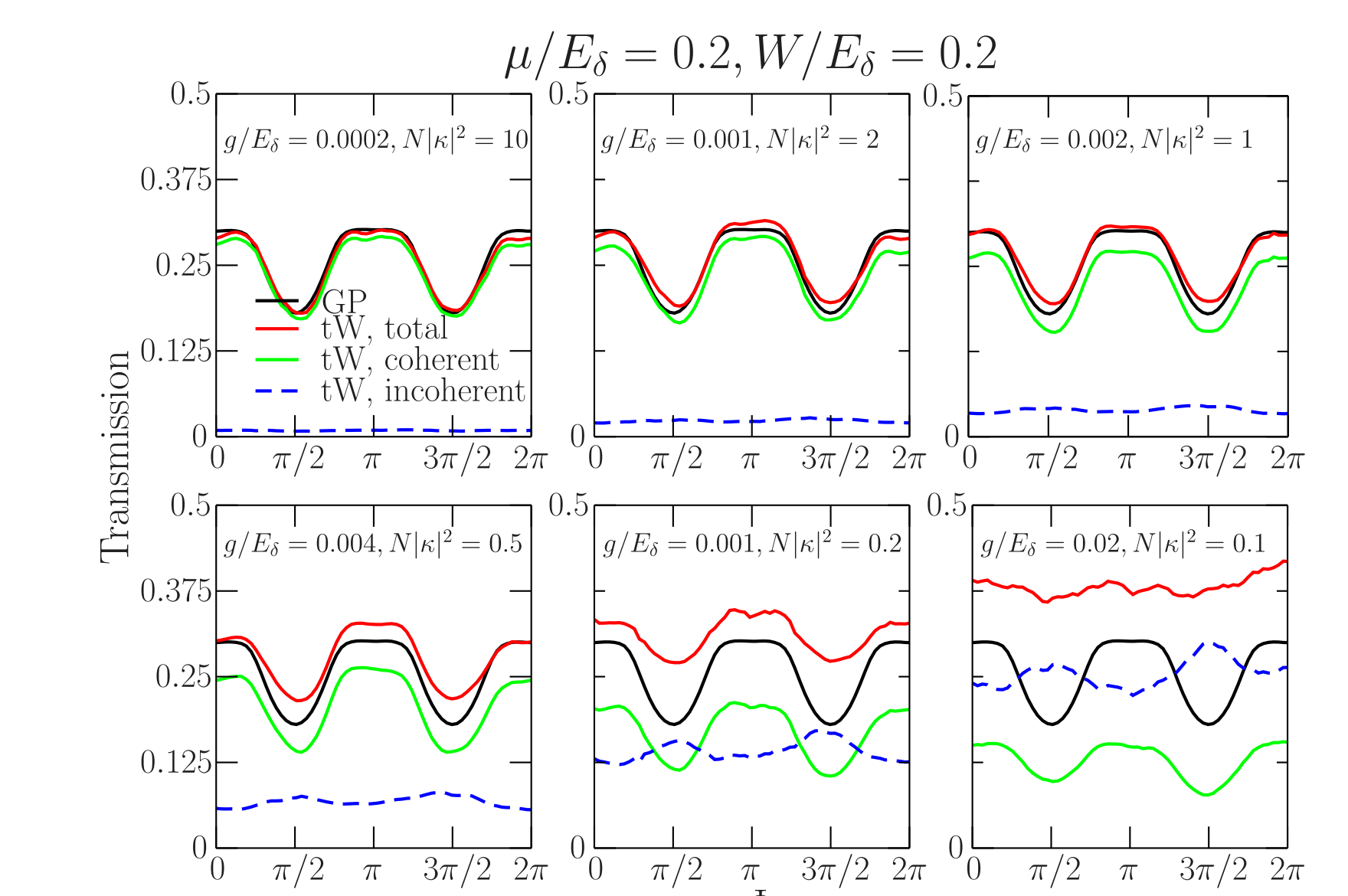


## AAS oscillations with interaction

What happens if we set a weak interaction ?



- The oscillations amplitude is reduced
- The minimum at  $\Phi = \pi$  becomes a maximum !



- Truncated Wigner simulations confirm the coherent peak inversion for weak interaction
- Presence of dephasing for strong interaction
- Analytical calculations with our 1D model more feasible
- Full diagrammatic theory with interaction (non-linearity)

[T. Hartmann et al. *Ann. Phys. (Amsterdam)* **327** (2012)]